## Topologicall cascade through vortex reconnection

## Renzo L. Ricca

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## Motivations

- Determine relationships between structural complexity of physical knots and energy.
- Quantify energy/helicity transfers in dynamical systems.


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- Extend and apply new topological techniques to study complex systems.


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Fluid structures in classical turbulence
Werle, ONERA 1974
(Van Dyke 1982)


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Vortex knots as tubular embeddings
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Physical embedding:

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by a standard foliation $\mathcal{F}_{\left\{p_{i}, q_{i}\right\}}$ of the $\boldsymbol{\omega}$-lines, such that $\boldsymbol{\omega} \cdot \hat{\boldsymbol{\nu}}=0$ on $\partial T$.


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- Definition: A vortex tangle $\mathcal{T}^{\prime}$ is a smooth immersion in $\mathbb{R}^{3}$ of finitely many disjoint standard tubular knots $K_{i}$, such that

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- Ideal evolution: circulation and topology preserved

$$
\Gamma_{i}=\int_{S_{i}} \omega \cdot \hat{\boldsymbol{\lambda}} d S=\text { cst. } ; \text { knot type } K_{i} \text { conserved. }
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Kinetic helicity and linking numbers

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H=\sum_{i}\left(\Gamma_{i}^{2} S L_{i}+\sum_{j \neq i} \Gamma_{i} \Gamma_{j} L k_{i j}\right) \quad\left\{\begin{array}{l}
S L_{i}=S L\left(K_{i}\right) \\
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## Topological cascade from head-on collision of vortex rings



(Klecker \& Irvine, Nature Physics 9 2013)

(Alekseenko et al., JETP Letters 7 2016)

## Mechanics of vortex reconnection

- Biot-Savart induction law:

$$
\mathbf{u}_{\mathrm{R}}(\mathrm{x})=\frac{1}{4 \pi} \int_{\Omega} \frac{\boldsymbol{\omega}\left(\mathrm{x}^{*}\right) \times\left(\mathrm{x}-\mathrm{x}^{*}\right)}{\left|\mathrm{x}-\mathrm{x}^{*}\right|^{3}} \mathrm{~d} V^{*}
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(Josserand \& Rossi 2007)


## Mechanics of vortex reconnection: pre-reconnection stage

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\mathbf{u}_{P}=\mathbf{u}_{\theta 1}+\mathbf{u}_{\theta 2}=\frac{\Gamma}{2 \pi \rho} \hat{\mathbf{e}}_{\theta 1}+\frac{\Gamma}{2 \pi \rho} \hat{\mathbf{e}}_{\theta 2}
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apply Bernoulli theorem at $P$ : $\quad p_{P}+\frac{1}{2}\left|\mathbf{u}_{P}\right|^{2}=$ constant

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\left|\mathbf{u}_{P}\right|^{2}=\left|\mathbf{u}_{\theta 1}+\mathbf{u}_{\theta 2}\right|^{2}=2\left(\frac{\Gamma}{2 \pi \rho}\right)^{2}(1+\cos \alpha): \quad\left|\mathbf{u}_{P}\right|_{\max }^{2} \Longleftrightarrow \alpha=0
$$

Observation of anti-parallel reconnection

(Alekseenko et al., JETP Letters 7 2016)

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$t^{*}=0$

$t^{*}=1.61$

$t^{*}=2.07$

$t^{*}=2.77$

DNS of vorticity iso-surfaces at 40\% of maximum initial vorticity (Hussain \& Duraisamy, Phys Fluids 23 2011)

Writhe conservation and helicity change under reconnection
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Corollary 2. Since under anti-parallel reconnection of $A$ and $B$ total torsion is conserved, any change in helicity is solely due to a change in intrinsic twist.

Vortex defects in condensates and superfluids


Taylor \& Dennis
(Nature Comm 2016)

Vortex defects in condensates and superfluids


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(c)


Cooper et al. (Nature Sci Rep 2019)

Tackling structural complexity by knot polynomials

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(Maxwell 1867)

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Knot polynomials?


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HOMFLYPT polynomial from self-linking

- Theorem (Liu \& Ricca, JFM 2015). If $K$ denotes a vortex knot of axis $C$ and helicity $H=H(K)$, then

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e^{H(K)}=e^{\oint_{C} \mathbf{u} \cdot d \mathbf{l}} \quad \text { appropriately rescaled }
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- Reduction of HOMFLYPT $P_{K}(a, z)$ to Jones $V_{K}(t)$ :
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a P(\nless)-a^{-1} P(\nearrow)={ }_{z} P()() \Longleftrightarrow[f(T w)]=g(W r) .
$$

- Reduction of HOMFLYPT $P_{K}(a, z)$ to Jones $V_{K}(t)$ :
$t=h(a, z): W r \propto T w \Longleftrightarrow$ knots and links are "framed".

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| $\boldsymbol{T}^{R}$ | $2 a^{-2}+a^{-2} z^{2}-a^{-4}$ | 1.51 |
| $\boldsymbol{F}^{8}$ | $a^{-2}-1-z^{2}+a^{2}$ | 0.17 |
| $\boldsymbol{W}$ | $\ldots a^{-1} z^{-1}-a^{-1} z+a z^{-1}+2 a z+a z^{3}-a^{3} z$ | 1.59 |
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\begin{aligned}
& \cdots+8 \rightarrow 8 \rightarrow \infty \rightarrow \infty \rightarrow \infty \\
& T(2,7) \quad T(2,6) \quad T(2,5) \quad T(2,4) \quad T(2,3) \quad T(2,2) \quad T(2,1) \quad T(2,0) \\
& \{T(2, n)\}: \ldots \rightarrow T(2,2 n+1) \rightarrow T(2, n) \rightarrow \ldots \rightarrow T(2,0) .
\end{aligned}
$$

- Theorem (Liu \& Ricca 2016). HOMFLYPT computation of $P_{T(2, n)}$ generates for decreasing $n$ a monotonically decreasing sequence of numerical values given by

$$
P_{T(2,3+q)}=A_{q}(\tau, \omega) P_{T(2,3)}+B_{q}(\tau, \omega) P_{T(2,2)} \quad(q \in \mathbb{N})
$$

where $A_{q}(\tau, \omega)$ and $B_{q}(\tau, \omega)$ are known functions of $\tau$ and $\omega$, with initial conditions $P_{T(2,3)}$ and $P_{T(2,2)} \cdot$

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| HOMFLYPT: <br> $a=e^{1 / 4}, k=e^{1 / 2}$ | 8.52 | 6.64 | 5.17 | 4.03 | 3.13 | 2.44 | 1.89 | 1.50 | 1.11 | 1 | 0.48 |  |  |  |  |  |  |
| Jones: <br> $\tau=e^{-1}$ | -0.01 | 0.02 | -0.03 | 0.05 | -0.09 | 0.15 | -0.25 | 0.40 | -0.69 | 1 | -2.26 |  |  |  |  |  |  |
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P_{T(2,4)}=3.13{ }^{P_{T(2,5)}=2.44 \rightarrow P_{T(2,4)}=1.89}{ }_{P_{T(2,3)}=1.50}
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## Selected references

- Liu, X. \& Ricca, R.L. (2012) The Jones polynomial for fluid knots from helicity. J Phys A: Math \& Theor 45, 205501.
- Laing, C.E., Ricca, R.L. \& Sumners, DeW. L. (2015) Conservation of writhe helicity under anti-parallel reconnection. Nature Scientific Reports 5, 9224.
- Liu, X. \& Ricca, R.L. (2015) On the derivation of HOMFLYPT polynomial invariant for fluid knots. J Fluid Mech 773, 34-48.
- Liu, X. \& Ricca, R.L. (2016) Knots cascade detected by a monotonically decreasing sequence of values. Nature Scientific Reports 6, 24118.
- Zuccher, S. \& Ricca, R.L. (2017) Relaxation of twist helicity in the cascade process of linked quantum vortices. Phys Rev E 95, 053109.
- Ricca, R.L. \& Liu, X. (2018) HOMFLYPT polynomial is the best quantifier for topological cascades of vortex knots. Fluid Dyn. Research 50, 011404.
- Oberti, C. \& Ricca, R.L. (2019) Influence of winding number on vortex knots dynamics. Nature Scientific Reports 9, 17284.

Thank you!

